

http://en.wikipedia.org/wiki/Fourier_transform

Operations

Operation	$x(t)$	$X(\omega)$			
Scalar multiplication	$kx(t)$	$kX(\omega)$			
Addition	$x_1(t) + x_2(t)$	$X_1(\omega) + X_2(\omega)$	Time convolution	$x_1(t) * x_2(t)$	$X_1(\omega)X_2(\omega)$
Conjugation	$x^*(t)$	$X^*(-\omega)$	Frequency convolution	$x_1(t)x_2(t)$	$\frac{1}{2\pi} X_1(\omega) * X_2(\omega)$
Duality	$X(t)$	$2\pi x(-\omega)$	Time differentiation	$\frac{d^n x}{dt^n}$	$(j\omega)^n X(\omega)$
Scaling (a real)	$x(at)$	$\frac{1}{ a } X\left(\frac{\omega}{a}\right)$	Time integration	$\int_{-\infty}^t x(u) du$	$\frac{X(\omega)}{j\omega} + \pi X(0)\delta(\omega)$
Time shifting	$x(t - t_0)$	$X(\omega)e^{-j\omega t_0}$			
Frequency shifting (ω_0 real)	$x(t)e^{j\omega_0 t}$	$X(\omega - \omega_0)$			

Functional relationships

	Function	Fourier transform unitary, ordinary frequency	Fourier transform unitary, angular frequency	Fourier transform non-unitary, angular frequency	Remarks
	$f(x)$	$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x)e^{-2\pi i x \xi} dx$	$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x)e^{-i\omega x} dx$	$\hat{f}(\nu) = \int_{-\infty}^{\infty} f(x)e^{-i\nu x} dx$	Definition
101	$a \cdot f(x) + b \cdot g(x)$	$a \cdot \hat{f}(\xi) + b \cdot \hat{g}(\xi)$	$a \cdot \hat{f}(\omega) + b \cdot \hat{g}(\omega)$	$a \cdot \hat{f}(\nu) + b \cdot \hat{g}(\nu)$	Linearity
102	$f(x - a)$	$e^{-2\pi i a \xi} \hat{f}(\xi)$	$e^{-i a \omega} \hat{f}(\omega)$	$e^{-i a \nu} \hat{f}(\nu)$	Shift in time domain
103	$e^{2\pi i a x} f(x)$	$\hat{f}(\xi - a)$	$\hat{f}(\omega - 2\pi a)$	$\hat{f}(\nu - 2\pi a)$	Shift in frequency domain, dual of 102
104	$f(ax)$	$\frac{1}{ a } \hat{f}\left(\frac{\xi}{a}\right)$	$\frac{1}{ a } \hat{f}\left(\frac{\omega}{a}\right)$	$\frac{1}{ a } \hat{f}\left(\frac{\nu}{a}\right)$	Scaling in the time domain. If $ a $ is large, then $f(ax)$ is concentrated around 0 and $\frac{1}{ a } \hat{f}\left(\frac{\omega}{a}\right)$ spreads out and flattens.
105	$\hat{f}(x)$	$f(-\xi)$	$f(-\omega)$	$2\pi f(-\nu)$	Duality. Here \hat{f} needs to be calculated using the same method as Fourier transform column. Results from swapping "dummy" variables of x and ξ or ω or ν .
106	$\frac{d^n f(x)}{dx^n}$	$(2\pi i \xi)^n \hat{f}(\xi)$	$(i\omega)^n \hat{f}(\omega)$	$(i\nu)^n \hat{f}(\nu)$	
107	$x^n f(x)$	$\left(\frac{i}{2\pi}\right)^n \frac{d^n \hat{f}(\xi)}{d\xi^n}$	$i^n \frac{d^n \hat{f}(\omega)}{d\omega^n}$	$i^n \frac{d^n \hat{f}(\nu)}{d\nu^n}$	This is the dual of 106
108	$(f * g)(x)$	$\hat{f}(\xi)\hat{g}(\xi)$	$\sqrt{2\pi} \hat{f}(\omega)\hat{g}(\omega)$	$\hat{f}(\nu)\hat{g}(\nu)$	The notation $f * g$ denotes the convolution of f and g — this rule is the convolution theorem
109	$f(x)g(x)$	$(\hat{f} * \hat{g})(\xi)$	$\frac{(\hat{f} * \hat{g})(\omega)}{\sqrt{2\pi}}$	$\frac{1}{2\pi} (\hat{f} * \hat{g})(\nu)$	This is the dual of 108
110	For $f(x)$ purely real	$\hat{f}(-\xi) = \overline{\hat{f}(\xi)}$	$\hat{f}(-\omega) = \overline{\hat{f}(\omega)}$	$\hat{f}(-\nu) = \overline{\hat{f}(\nu)}$	Hermitian symmetry. \bar{z} indicates the complex conjugate .
111	For $f(x)$ a purely real even function	$\hat{f}(\omega)$, $\hat{f}(\xi)$ and $\hat{f}(\nu)$ are purely real even functions .			
112	For $f(x)$ a purely real odd function	$\hat{f}(\omega)$, $\hat{f}(\xi)$ and $\hat{f}(\nu)$ are purely imaginary odd functions .			

Square-integrable functions

	Function	Fourier transform unitary, ordinary frequency	Fourier transform unitary, angular frequency	Fourier transform non-unitary, angular frequency	Remarks
	$f(x)$	$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x)e^{-2\pi i x \xi} dx$	$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x)e^{-i\omega x} dx$	$\hat{f}(\nu) = \int_{-\infty}^{\infty} f(x)e^{-i\nu x} dx$	
201	$\text{rect}(ax)$	$\frac{1}{ a } \cdot \text{sinc}\left(\frac{\xi}{a}\right)$	$\frac{1}{\sqrt{2\pi a^2}} \cdot \text{sinc}\left(\frac{\omega}{2\pi a}\right)$	$\frac{1}{ a } \cdot \text{sinc}\left(\frac{\nu}{2\pi a}\right)$	The rectangular pulse and the normalized sinc function , here defined as $\text{sinc}(x) = \sin(\pi x)/(\pi x)$
202	$\text{sinc}(ax)$	$\frac{1}{ a } \cdot \text{rect}\left(\frac{\xi}{a}\right)$	$\frac{1}{\sqrt{2\pi a^2}} \cdot \text{rect}\left(\frac{\omega}{2\pi a}\right)$	$\frac{1}{ a } \cdot \text{rect}\left(\frac{\nu}{2\pi a}\right)$	Dual of rule 201. The rectangular function is an ideal low-pass filter , and the sinc function is the non-causal impulse response of such a filter.
203	$\text{sinc}^2(ax)$	$\frac{1}{ a } \cdot \text{tri}\left(\frac{\xi}{a}\right)$	$\frac{1}{\sqrt{2\pi a^2}} \cdot \text{tri}\left(\frac{\omega}{2\pi a}\right)$	$\frac{1}{ a } \cdot \text{tri}\left(\frac{\nu}{2\pi a}\right)$	The function $\text{tri}(x)$ is the triangular function
204	$\text{tri}(ax)$	$\frac{1}{ a } \cdot \text{sinc}^2\left(\frac{\xi}{a}\right)$	$\frac{1}{\sqrt{2\pi a^2}} \cdot \text{sinc}^2\left(\frac{\omega}{2\pi a}\right)$	$\frac{1}{ a } \cdot \text{sinc}^2\left(\frac{\nu}{2\pi a}\right)$	Dual of rule 203.
205	$e^{-ax}u(x)$	$\frac{1}{a + 2\pi i \xi}$	$\frac{1}{\sqrt{2\pi}(a + i\omega)}$	$\frac{1}{a + i\nu}$	The function $u(x)$ is the Heaviside unit step function and $a > 0$.
206	$e^{-\alpha x^2}$	$\sqrt{\frac{\pi}{\alpha}} \cdot e^{-\frac{(\pi\xi)^2}{\alpha}}$	$\frac{1}{\sqrt{2\alpha}} \cdot e^{-\frac{\omega^2}{4\alpha}}$	$\sqrt{\frac{\pi}{\alpha}} \cdot e^{-\frac{\nu^2}{4\alpha}}$	This shows that, for the unitary Fourier transforms, the Gaussian function $\exp(-ax^2)$ is its own Fourier transform for some choice of α . For this to be integrable we must have $\text{Re}(\alpha) > 0$.
207	$e^{-a x }$	$\frac{2a}{a^2 + 4\pi^2\xi^2}$	$\sqrt{\frac{2}{\pi}} \cdot \frac{a}{a^2 + \omega^2}$	$\frac{2a}{a^2 + \nu^2}$	For $a > 0$. That is, the Fourier transform of a decaying exponential function is a Lorentzian function .
208	$\text{sech}(ax)$	$\frac{\pi}{a} \text{sech}\left(\frac{\pi^2}{a}\xi\right)$	$\frac{1}{a} \sqrt{\frac{\pi}{2}} \text{sech}\left(\frac{\pi}{2a}\omega\right)$	$\frac{\pi}{a} \text{sech}\left(\frac{\pi}{2a}\nu\right)$	Hyperbolic secant is its own Fourier transform
209	$e^{-\frac{a^2 x^2}{2}} H_n(ax)$	$\frac{\sqrt{2\pi}(-i)^n}{a} \cdot e^{-\frac{2\pi^2\xi^2}{a^2}} H_n\left(\frac{2\pi\xi}{a}\right)$	$\frac{(-i)^n}{a} \cdot e^{-\frac{\omega^2}{2a^2}} H_n\left(\frac{\omega}{a}\right)$	$\frac{(-i)^n \sqrt{2\pi}}{a} \cdot e^{-\frac{\nu^2}{2a^2}} H_n\left(\frac{\nu}{a}\right)$	H_n is the Hermite's polynomial . If $a = 1$ then the Gauss-Hermite functions are eigenfunctions of the Fourier transform operator. For a derivation, see Hermite polynomial . The formula reduces to 206 for $n = 0$.

Distributions

	Function	Fourier transform unitary, ordinary frequency	Fourier transform unitary, angular frequency	Fourier transform non-unitary, angular frequency	Remarks
	$f(x)$	$\hat{f}(\xi) = \int_{-\infty}^{\infty} f(x)e^{-2\pi i x \xi} dx$	$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x)e^{-i\omega x} dx$	$\hat{f}(\nu) = \int_{-\infty}^{\infty} f(x)e^{-i\nu x} dx$	
301	1	$\delta(\xi)$	$\sqrt{2\pi} \cdot \delta(\omega)$	$2\pi\delta(\nu)$	The distribution $\delta(\xi)$ denotes the Dirac delta function .
302	$\delta(x)$	1	$\frac{1}{\sqrt{2\pi}}$	1	Dual of rule 301.
303	e^{iax}	$\delta\left(\xi - \frac{a}{2\pi}\right)$	$\sqrt{2\pi} \cdot \delta(\omega - a)$	$2\pi\delta(\nu - a)$	This follows from 103 and 301.
304	$\cos(ax)$	$\frac{\delta\left(\xi - \frac{a}{2\pi}\right) + \delta\left(\xi + \frac{a}{2\pi}\right)}{2}$	$\sqrt{2\pi} \cdot \frac{\delta(\omega - a) + \delta(\omega + a)}{2}$	$\pi(\delta(\nu - a) + \delta(\nu + a))$	This follows from rules 101 and 303 using Euler's formula : $\cos(ax) = (e^{iax} + e^{-iax})/2$.
305	$\sin(ax)$	$\frac{\delta\left(\xi - \frac{a}{2\pi}\right) - \delta\left(\xi + \frac{a}{2\pi}\right)}{2i}$	$\sqrt{2\pi} \cdot \frac{\delta(\omega - a) - \delta(\omega + a)}{2i}$	$-i\pi(\delta(\nu - a) - \delta(\nu + a))$	This follows from 101 and 303 using $\sin(ax) = (e^{iax} - e^{-iax})/(2i)$.
306	$\cos(ax^2)$	$\sqrt{\frac{\pi}{a}} \cos\left(\frac{\pi^2 \xi^2}{a} - \frac{\pi}{4}\right)$	$\frac{1}{\sqrt{2a}} \cos\left(\frac{\omega^2}{4a} - \frac{\pi}{4}\right)$	$\sqrt{\frac{\pi}{a}} \cos\left(\frac{\nu^2}{4a} - \frac{\pi}{4}\right)$	
307	$\sin(ax^2)$	$-\sqrt{\frac{\pi}{a}} \sin\left(\frac{\pi^2 \xi^2}{a} - \frac{\pi}{4}\right)$	$\frac{-1}{\sqrt{2a}} \sin\left(\frac{\omega^2}{4a} - \frac{\pi}{4}\right)$	$-\sqrt{\frac{\pi}{a}} \sin\left(\frac{\nu^2}{4a} - \frac{\pi}{4}\right)$	
308	x^n	$\left(\frac{i}{2\pi}\right)^n \delta^{(n)}(\xi)$	$i^n \sqrt{2\pi} \delta^{(n)}(\omega)$	$2\pi i^n \delta^{(n)}(\nu)$	Here, n is a natural number and $\delta^{(n)}(\xi)$ is the n -th distribution derivative of the Dirac delta function. This rule follows from rules 107 and 301. Combining this rule with 101, we can transform all polynomials .
309	$\frac{1}{x}$	$-i\pi \operatorname{sgn}(\xi)$	$-i\sqrt{\frac{\pi}{2}} \operatorname{sgn}(\omega)$	$-i\pi \operatorname{sgn}(\nu)$	Here $\operatorname{sgn}(\xi)$ is the sign function . Note that $1/x$ is not a distribution. It is necessary to use the Cauchy principal value when testing against Schwartz functions. This rule is useful in studying the Hilbert transform .
310	$\frac{1}{x^n} := \frac{(-1)^{n-1}}{(n-1)!} \frac{d^{n-1}}{dx^{n-1}} \log x $	$-i\pi \frac{(-2\pi i \xi)^{n-1}}{(n-1)!} \operatorname{sgn}(\xi)$	$-i\sqrt{\frac{\pi}{2}} \cdot \frac{(-i\omega)^{n-1}}{(n-1)!} \operatorname{sgn}(\omega)$	$-i\pi \frac{(-i\nu)^{n-1}}{(n-1)!} \operatorname{sgn}(\nu)$	$1/x^n$ is the homogeneous distribution defined by the distributional derivative $\frac{(-1)^{n-1}}{(n-1)!} \frac{d^{n-1}}{dx^{n-1}} \log x $

311	$ x ^\alpha$	$-2 \frac{\sin(\pi\alpha/2)\Gamma(\alpha+1)}{ 2\pi\xi ^{\alpha+1}}$	$\frac{-2}{\sqrt{2\pi}} \frac{\sin(\pi\alpha/2)\Gamma(\alpha+1)}{ \omega ^{\alpha+1}}$	$-2 \frac{\sin(\pi\alpha/2)\Gamma(\alpha+1)}{ \nu ^{\alpha+1}}$	This formula is valid for $0 > \alpha > -1$. For $\alpha > 0$ some singular terms arise at the origin that can be found by differentiating 318. If $\text{Re } \alpha > -1$, then $ x ^\alpha$ is a locally integrable function, and so a tempered distribution. The function $\alpha \mapsto x ^\alpha$ is a holomorphic function from the right half-plane to the space of tempered distributions. It admits a unique meromorphic extension to a tempered distribution, also denoted $ x ^\alpha$ for $\alpha \neq -2, -4, \dots$ (See homogeneous distribution .)
312	$\text{sgn}(x)$	$\frac{1}{i\pi\xi}$	$\sqrt{\frac{2}{\pi}} \frac{1}{i\omega}$	$\frac{2}{i\nu}$	The dual of rule 309. This time the Fourier transforms need to be considered as Cauchy principal value .
313	$u(x)$	$\frac{1}{2} \left(\frac{1}{i\pi\xi} + \delta(\xi) \right)$	$\sqrt{\frac{\pi}{2}} \left(\frac{1}{i\pi\omega} + \delta(\omega) \right)$	$\pi \left(\frac{1}{i\pi\nu} + \delta(\nu) \right)$	The function $u(x)$ is the Heaviside unit step function ; this follows from rules 101, 301, and 312.
314	$\sum_{n=-\infty}^{\infty} \delta(x - nT)$	$\frac{1}{T} \sum_{k=-\infty}^{\infty} \delta\left(\xi - \frac{k}{T}\right)$	$\frac{\sqrt{2\pi}}{T} \sum_{k=-\infty}^{\infty} \delta\left(\omega - \frac{2\pi k}{T}\right)$	$\frac{2\pi}{T} \sum_{k=-\infty}^{\infty} \delta\left(\nu - \frac{2\pi k}{T}\right)$	This function is known as the Dirac comb function . This result can be derived from 302 and 102, together with the fact that $\sum_{n=-\infty}^{\infty} e^{inx} = 2\pi \sum_{k=-\infty}^{\infty} \delta(x + 2\pi k)$ as distributions.
315	$J_0(x)$	$\frac{2 \text{rect}(\pi\xi)}{\sqrt{1 - 4\pi^2\xi^2}}$	$\sqrt{\frac{2}{\pi}} \cdot \frac{\text{rect}\left(\frac{\omega}{2}\right)}{\sqrt{1 - \omega^2}}$	$\frac{2 \text{rect}\left(\frac{\nu}{2}\right)}{\sqrt{1 - \nu^2}}$	The function $J_0(x)$ is the zeroth order Bessel function of first kind.
316	$J_n(x)$	$\frac{2(-i)^n T_n(2\pi\xi) \text{rect}(\pi\xi)}{\sqrt{1 - 4\pi^2\xi^2}}$	$\sqrt{\frac{2}{\pi}} \frac{(-i)^n T_n(\omega) \text{rect}\left(\frac{\omega}{2}\right)}{\sqrt{1 - \omega^2}}$	$\frac{2(-i)^n T_n(\nu) \text{rect}\left(\frac{\nu}{2}\right)}{\sqrt{1 - \nu^2}}$	This is a generalization of 315. The function $J_n(x)$ is the n -th order Bessel function of first kind. The function $T_n(x)$ is the Chebyshev polynomial of the first kind.
317	$\log x $	$-\frac{1}{2} \frac{1}{ \xi } - \gamma\delta(\xi)$	$-\frac{\sqrt{\pi/2}}{ \omega } - \sqrt{2\pi}\gamma\delta(\omega)$	$-\frac{\pi}{ \nu } - 2\pi\gamma\delta(\nu)$	γ is the Euler–Mascheroni constant .
318	$(\mp ix)^{-\alpha}$	$\frac{(2\pi)^\alpha}{\Gamma(\alpha)} u(\pm\xi) (\pm\xi)^{\alpha-1}$	$\frac{\sqrt{2\pi}}{\Gamma(\alpha)} u(\pm\omega) (\pm\omega)^{\alpha-1}$	$\frac{2\pi}{\Gamma(\alpha)} u(\pm\nu) (\pm\nu)^{\alpha-1}$	This formula is valid for $1 > \alpha > 0$. Use differentiation to derive formula for higher exponents. u is the Heaviside function.

Two-dimensional function

	Function	Fourier transform unitary, ordinary frequency	Fourier transform unitary, angular frequency	Fourier transform non-unitary, angular frequency
400	$f(x, y)$	$\hat{f}(\xi_x, \xi_y) = \iint f(x, y) e^{-2\pi i(\xi_x x + \xi_y y)} dx dy$	$\hat{f}(\omega_x, \omega_y) = \frac{1}{2\pi} \iint f(x, y) e^{-i(\omega_x x + \omega_y y)} dx dy$	$\hat{f}(\nu_x, \nu_y) = \iint f(x, y) e^{-i(\nu_x x + \nu_y y)} dx dy$
401	$e^{-\pi(a^2 x^2 + b^2 y^2)}$	$\frac{1}{ ab } e^{-\pi(\xi_x^2/a^2 + \xi_y^2/b^2)}$	$\frac{1}{2\pi \cdot ab } e^{-\frac{(\omega_x^2/a^2 + \omega_y^2/b^2)}{4\pi}}$	$\frac{1}{ ab } e^{-\frac{(\nu_x^2/a^2 + \nu_y^2/b^2)}{4\pi}}$
402	$\text{circ}(\sqrt{x^2 + y^2})$	$\frac{J_1(2\pi\sqrt{\xi_x^2 + \xi_y^2})}{\sqrt{\xi_x^2 + \xi_y^2}}$	$\frac{J_1(\sqrt{\omega_x^2 + \omega_y^2})}{\sqrt{\omega_x^2 + \omega_y^2}}$	$\frac{2\pi J_1(\sqrt{\nu_x^2 + \nu_y^2})}{\sqrt{\nu_x^2 + \nu_y^2}}$

Formulas for general n-dimensional functions

	Function	Fourier transform unitary, ordinary frequency	Fourier transform unitary, angular frequency	Fourier transform non-unitary, angular frequency
500	$f(x)$	$\hat{f}(\xi) = \int_{\mathbf{R}^n} f(x) e^{-2\pi i x \cdot \xi} d^n x$	$\hat{f}(\omega) = \frac{1}{(2\pi)^{(n/2)}} \int_{\mathbf{R}^n} f(x) e^{-i\omega \cdot x} d^n x$	$\hat{f}(\nu) = \int_{\mathbf{R}^n} f(x) e^{-i x \cdot \nu} d^n x$
501	$\chi_{[0,1]}(x)(1 - x ^2)^\delta$	$\pi^{-\delta} \Gamma(\delta + 1) \xi ^{-(n/2) - \delta} \cdot J_{n/2 + \delta}(2\pi \xi)$	$2^{-\delta} \Gamma(\delta + 1) \omega ^{-(n/2) - \delta} \cdot J_{n/2 + \delta}(\omega)$	$\pi^{-\delta} \Gamma(\delta + 1) \left \frac{\nu}{2\pi} \right ^{-(n/2) - \delta} \cdot J_{n/2 + \delta}(\nu)$
502	$ x ^{-\alpha}, \quad 0 < \text{Re } \alpha < n.$	$c_\alpha \xi ^{-(n-\alpha)}$		